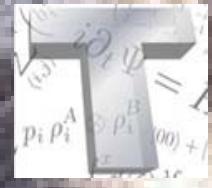




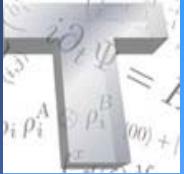
**COMPAS meeting**  
**November 30th – December 1st**  
**Hotel Le Dôme, Brussels**



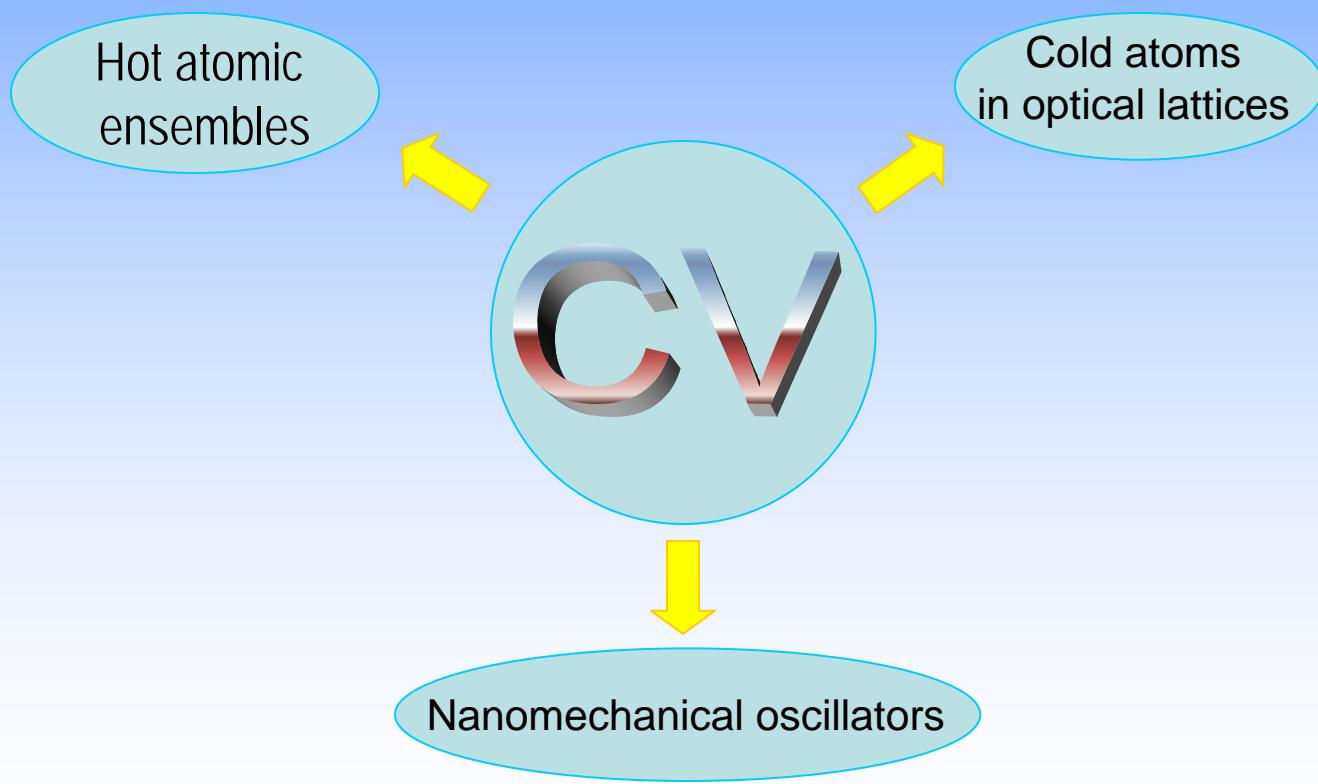
Christine Muschik

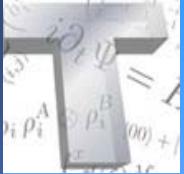


J.Ignacio Cirac  
Max-Planck-Institut für Quantenoptik



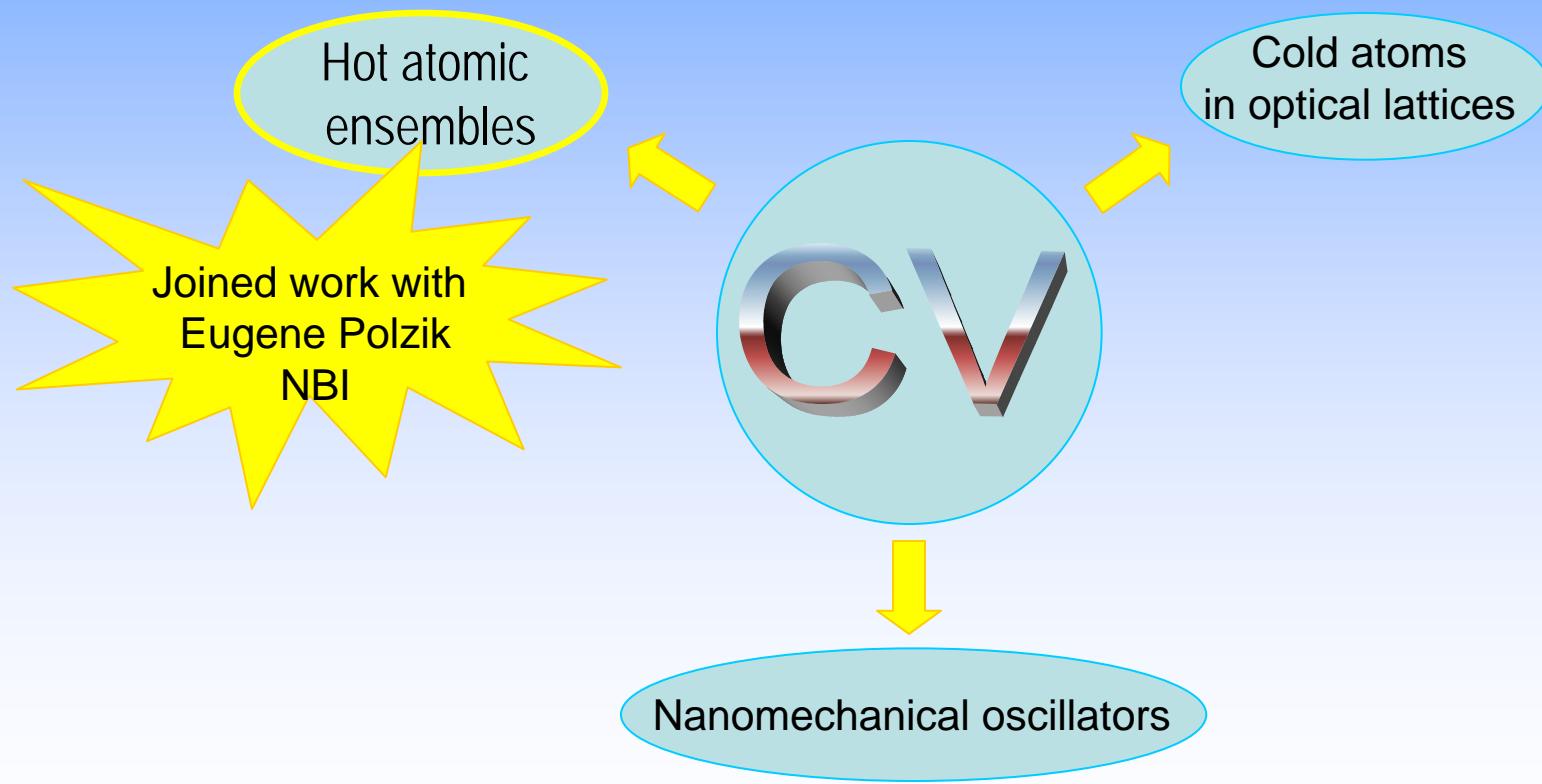
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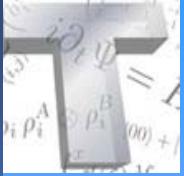




# COMPAS meeting

November 30th – December 1st  
Hotel Le Dôme, Brussels





COMPAS meeting  
November 30th – December 1st  
Hotel Le Dôme, Brussels

# Steady –state entanglement with atomic ensembles

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Christine Muschik and J.Ignacio Cirac

Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Strasse 1, 85748 Garching, Germany

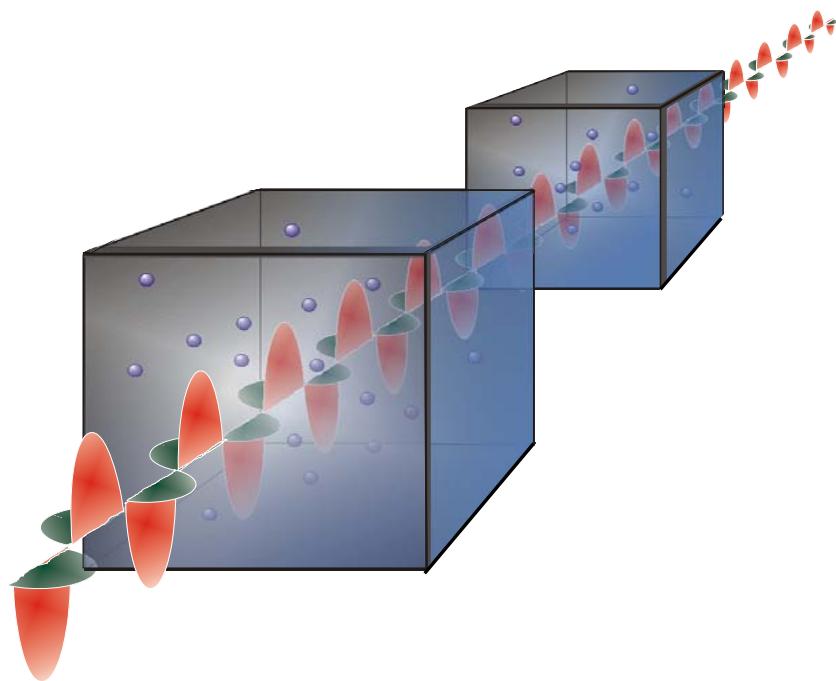
Hanna Krauter, Kasper Jensen, Wojciech Wasilewski,  
Jelmer Renema and Eugene Polzik

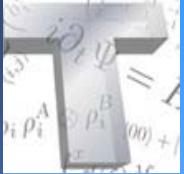
Niels Bohr Institute, Danish Research Foundation Center for Quantum Optics (QUANTOP),  
Blegdamsvej 17, DK-2100 Copenhagen, Denmark

## Overview:

### Steady state entanglement with atomic ensembles

- Motivation
- Main idea
- Theoretical model
- Experimental realization





## Motivation:

### Quantum entanglement

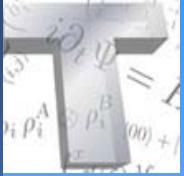
- Basic ingredient in most applications in the field of quantum information
- But: Lifetimes are usually very short
- Therefore: Need for much longer lifetimes!

### Typically:

- Quantum states are fragile under decoherence
- Avoidance of dissipation by decoupling the system from the environment
- Strict isolation!

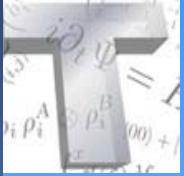
# ENVIRONMENT





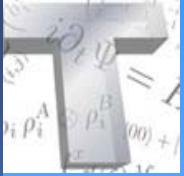
# Motivation:





## Motivation:





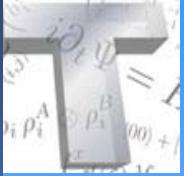
## Motivation:

New approaches

S. Diehl, A. Micheli, A.  
Kantian, B. Kraus, H.P.  
Büchler, P. Zoller  
  
F. Verstraete, M.M.  
Wolf, J.I. Cirac,

→ Use the interaction of the system with the environment





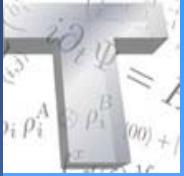
## Motivation:



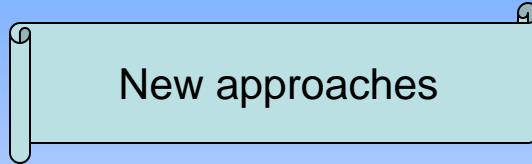
New approaches

- Use the interaction of the system with the environment
- Dissipation drives the system into the desired state



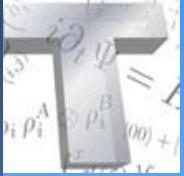


## Motivation:



- New approaches
- Use the interaction of the system with the environment
- Dissipation drives the system into the desired state
- Robust method to create extremely long-lived and robust entanglement



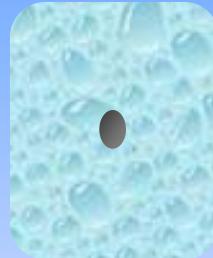


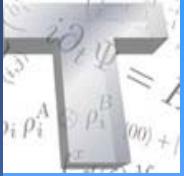
## Motivation:



### Procedure:

- Engineer the coupling
- Steady state = desired state
- Arbitrary initial state

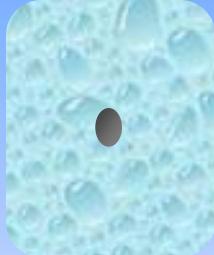




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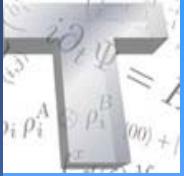
### Procedure:

- Engineer the coupling
- Steady state = desired state
- Arbitrary initial state



### Proposals:

- Single trapped ion  
J. F. Poyatos, J. I. Cirac and P. Zoller, Phys. Rev. Lett. **77**, 4728 (1996)
- Atoms in two cavities  
B. Kraus and J. I. Cirac, Phys. Rev. Lett. **92**, 013602 (2004)
- Many-body systems  
S. Diehl, A. Micheli, A. Kantian, B. Kraus, H.P. Büchler, P. Zoller, Nature Physics **4**, 878 (2008)  
F. Verstraete, M.M. Wolf, J.I. Cirac, Nature Physics **5**, 633 (2009)
  - Quantum phase transitions
  - State preparation
  - Quantum computing

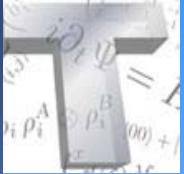


## Main idea:



Dissipative evolution:

$$d_t \rho(t) = \Gamma (2A\rho(t)A^\dagger - A^\dagger A \rho(t) - \rho(t)A^\dagger A)$$



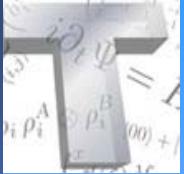
## Main idea:



Dissipative evolution:

$$d_t \rho(t) = \Gamma(2A\rho(t)A^\perp - A^\perp A \rho(t) - \rho(t)A^\perp A)$$

$A|\Psi\rangle = 0 \Rightarrow \rho_\Psi = |\Psi\rangle\langle\Psi|$  is the unique steady state



## Main idea:



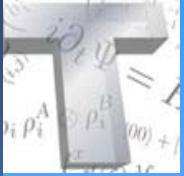
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Target state: Two mode squeezed state:  $\Psi_{EPR}$

$$A|\Psi_{EPR}\rangle = B|\Psi_{EPR}\rangle = 0$$



## Main idea:



Dissipative evolution:

$$d_t \rho(t) = \Gamma(2A\rho(t)A^+ - A^+A\rho(t) - \rho(t)A^+A)$$

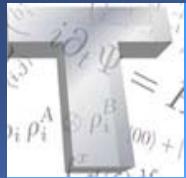
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**Target state:** Two mode squeezed state:  $\Psi_{EPR}$

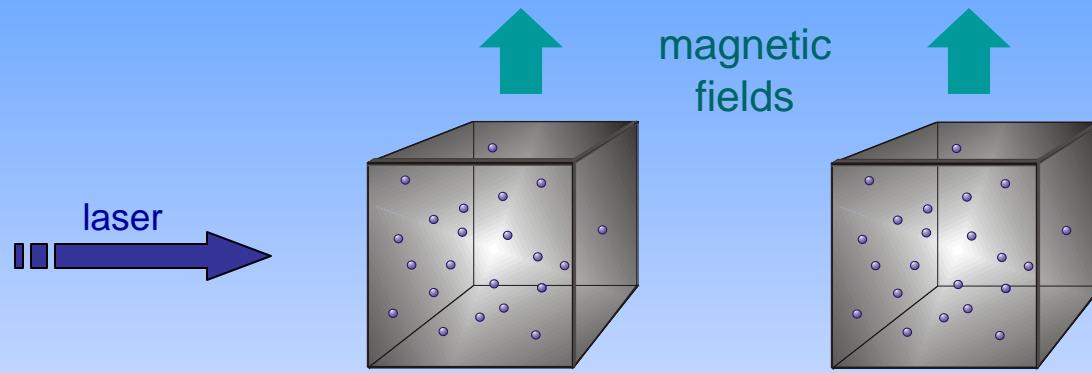
$$A|\Psi_{EPR}\rangle = B|\Psi_{EPR}\rangle = 0$$

$\rho_{EPR} = |\Psi_{EPR}\rangle\langle\Psi_{EPR}|$  is the unique steady state of the dissipative evolution governed by

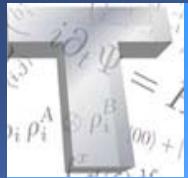
$$d_t \rho(t) = \Gamma(2A\rho(t)A^+ - A^+A\rho(t) - \rho(t)A^+A) + \Gamma(2B\rho(t)B^+ - B^+B\rho(t) - \rho(t)B^+B)$$



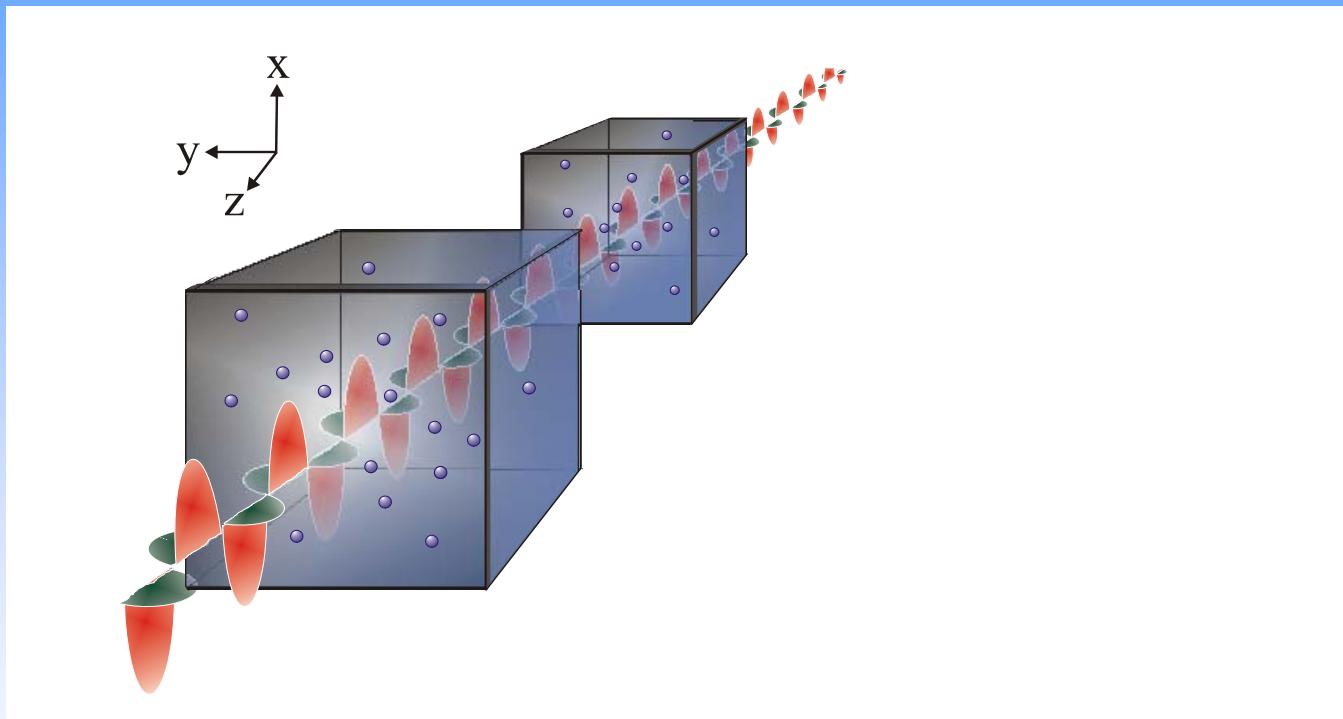
## Setup:

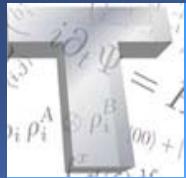


- Reservoir: common modes of the electromagnetic field.
- Control: Laser and magnetic fields

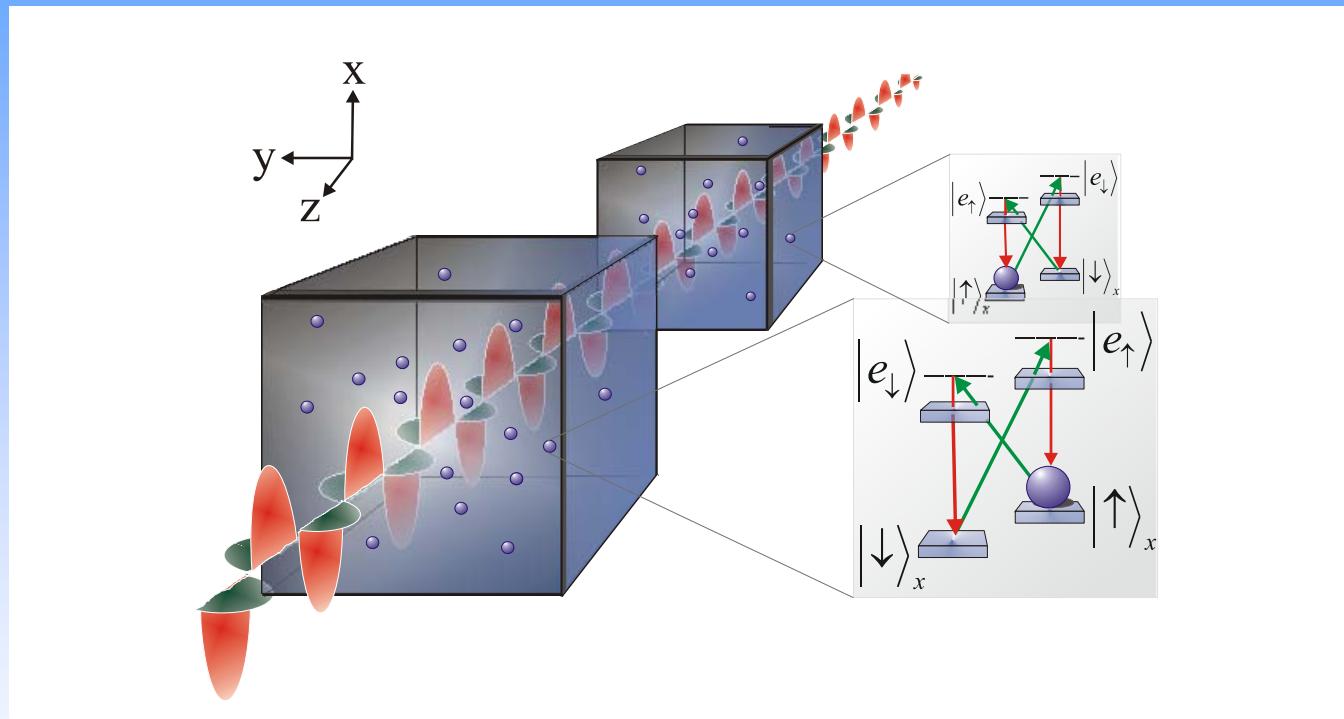


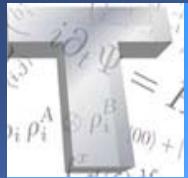
# Setup:





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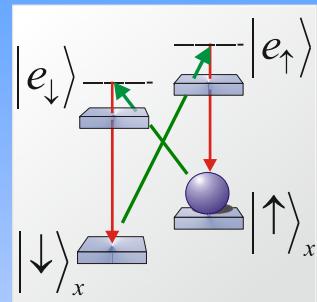




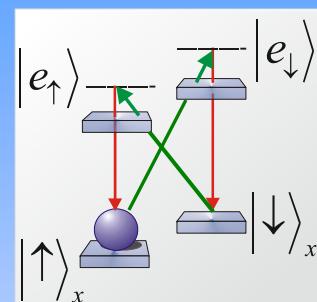
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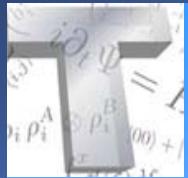


Ensemble 1



Ensemble 2

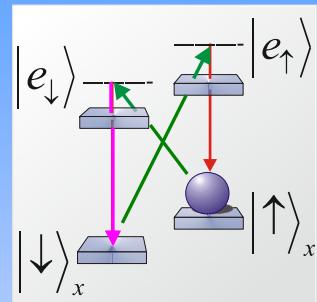




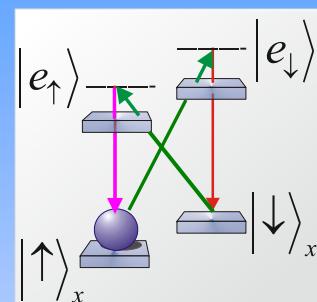
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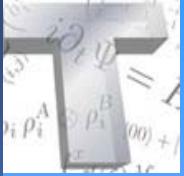


Ensemble 1



Ensemble 2

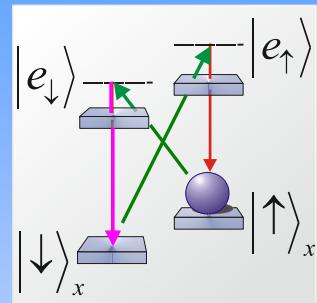




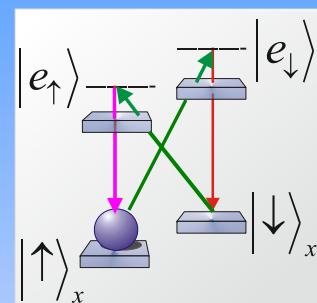
# Theory:



Ensemble 1

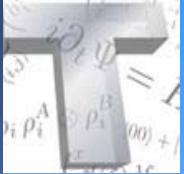


Ensemble 2



Processes:

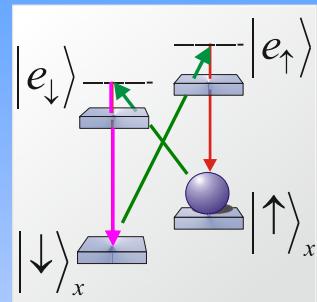
$$\text{Processes: } \downarrow (\mu\sigma_1^- + \nu\sigma_2^+)a_k \quad \downarrow (\nu\sigma_1^+ + \mu\sigma_2^-)a_k$$



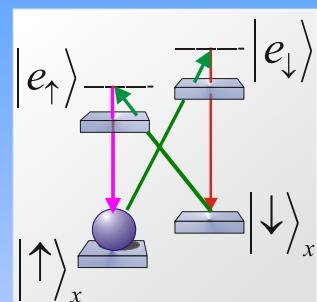
# Theory:



Ensemble 1



Ensemble 2



Processes:



$$A^+ a_k$$

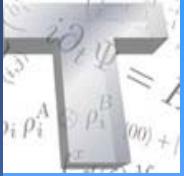


$$B^+ a_k$$

$$A = \mu \frac{1}{\sqrt{N}} \sum_i \sigma_{1,i}^+ + \nu \frac{1}{\sqrt{N}} \sum_j \sigma_{2,j}^-$$

$$\mu^2 - \nu^2 = 1$$

$$B = \nu \frac{1}{\sqrt{N}} \sum_i \sigma_{1,i}^- + \mu \frac{1}{\sqrt{N}} \sum_j \sigma_{2,j}^+$$

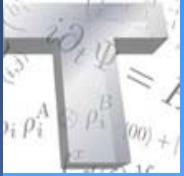


## Theory:



⇒ Masterequation:

$$\frac{d\rho}{dt} = \Gamma d(2A\rho A^\dagger - A^\dagger A \rho - \rho A^\dagger A) + \Gamma d(2B\rho B^\dagger - B^\dagger B \rho - \rho B^\dagger B)$$



## Theory:

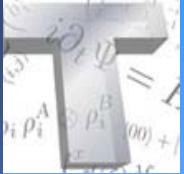


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+ undesired processes

- Adiabatic elimination excited states
- Two independent bands of modes
- Born-Markov approximation
- Room temperature (average atomic motion)



## Theory:



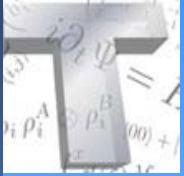
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+undesired processes

- Dark state:  $A |\Psi\rangle = 0$   
 $B |\Psi\rangle = 0$
- Entanglement: ideal case

$$\xi = \frac{\text{var}(J_{z,I} + J_{z,II}) + \text{var}(J_{y,I} - J_{y,II})}{\langle J_{x,I} \rangle + \langle J_{x,II} \rangle} = (\mu - \nu)^2$$

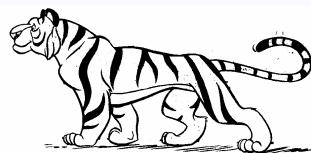


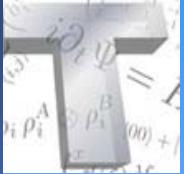
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## Theory:



- Entanglement: ideal case

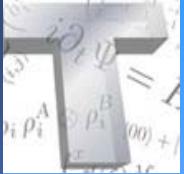
$$\xi = \frac{\text{var}(J_{z,I} + J_{z,II}) + \text{var}(J_{y,I} - J_{y,II})}{\langle J_{x,I} \rangle + \langle J_{x,II} \rangle} = (\mu - \nu)^2$$

- Entanglement: including undesired processes

$$\xi = \frac{1}{P_2} \frac{\tilde{\Gamma} + d\Gamma P_2^2 (\mu - \nu)^2}{\tilde{\Gamma} + d\Gamma P_2} \xrightarrow{d \gg 1} (\mu - \nu)^2$$

$\tilde{\Gamma}$  : noise rate

$P_2$  : polarization



## Theory:



- Entanglement: ideal case

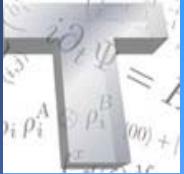
$$\xi = \frac{\text{var}(J_{z,I} + J_{z,II}) + \text{var}(J_{y,I} - J_{y,II})}{\langle J_{x,I} \rangle + \langle J_{x,II} \rangle} = (\mu - \nu)^2$$

- Entanglement: including undesired processes

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$\bar{\Gamma}$  : noise rate

$P_2$  : polarization



Effective ground state Hamiltonian:

$$H = \int_{\Delta\omega_{us}} d\vec{k} \cdot g(\omega_k) \left( \mu \sum_{i=1}^N \sigma_{I,i} e^{i\Delta\vec{k}\vec{r}_i} + \nu \sum_{j=1}^N \sigma_{II,j}^\perp e^{i\Delta\vec{k}\vec{r}_j} \right) a_{\vec{k}}^\perp + \int_{\Delta\omega_{ls}} d\vec{k} \cdot g(\omega_k) \left( \nu \sum_{i=1}^N \sigma_{I,i}^\perp e^{i\Delta\vec{k}\vec{r}_i} + \mu \sum_{j=1}^N \sigma_{II,j} e^{i\Delta\vec{k}\vec{r}_j} \right) a_{\vec{k}}^\perp + H.C.$$

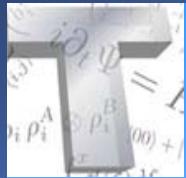
Master equation:

$$\begin{aligned} d_t \rho(t)_{first\ term} &= 2d\Gamma A \rho(t) A^\perp + 2d\Gamma B \rho(t) B^\perp \\ &+ 2\Gamma_{cool} \sum_{i=1}^N (\sigma_{I,i} \rho(t) \sigma_{I,i}^\perp + \sigma_{II,i} \rho(t) \sigma_{II,i}^\perp) \\ &+ 2\Gamma_{heat} \sum_{i=1}^N (\sigma_{I,i}^\perp \rho(t) \sigma_{I,i} + \sigma_{II,i}^\perp \rho(t) \sigma_{II,i}) \\ &+ 2\Gamma_{deph} \sum_{i=1}^N (\sigma_{I,i} \sigma_{I,i}^\perp \rho(t) \sigma_{I,i} \sigma_{I,i}^\perp + \sigma_{II,i} \sigma_{II,i}^\perp \rho(t) \sigma_{II,i} \sigma_{II,i}^\perp) \end{aligned}$$

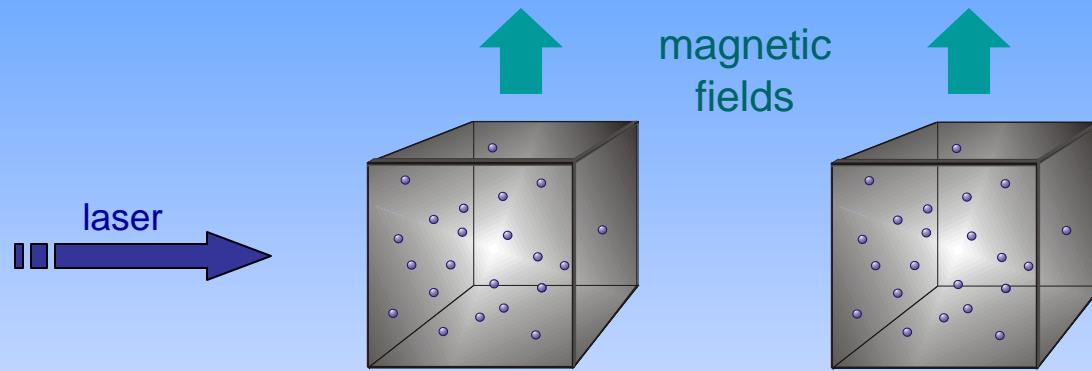
Entanglement:

$$\xi(t) = \frac{e^{-2(\tilde{\Gamma} + d\Gamma P_2(t))t}}{P_2(t)} + \frac{n_2(t)}{P_2(t)} \frac{\tilde{\Gamma} + d\Gamma P_2^2(t)(\mu - \nu)^2}{\tilde{\Gamma} + d\Gamma P_2(t)} \left( 1 - e^{-2(\tilde{\Gamma} + d\Gamma P_2(t))t} \right)$$

$$\tilde{\Gamma} = \Gamma_{cool} + \Gamma_{heat} + \Gamma_{deph}$$

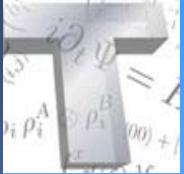


## Theory:

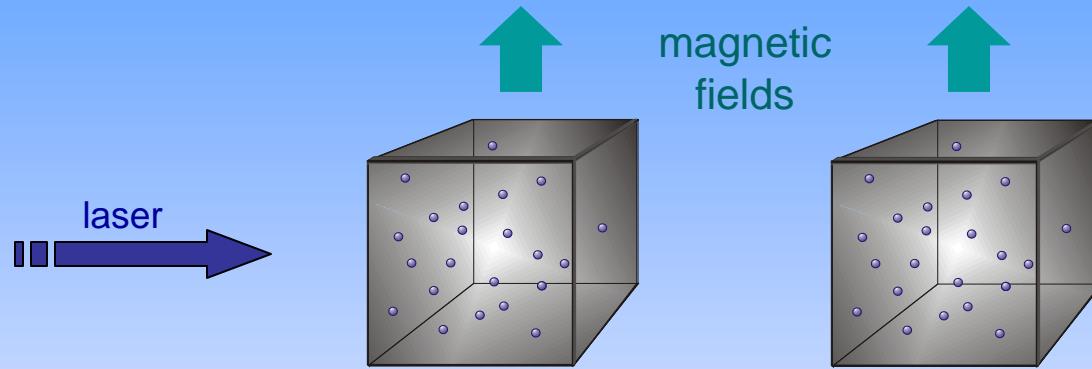


→ Result:

- Steady-state is an entangled state
- Immune to noise
- Long-lived entanglement



## Theory:

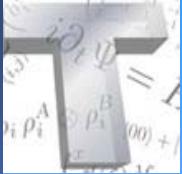


→ Result:

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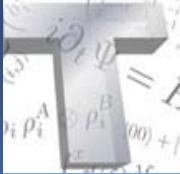
→ “Standard procedure“:

- Pure state (cooling, polarization, etc.)
- Coherent interaction (gate)
- Isolation (no decoherence)

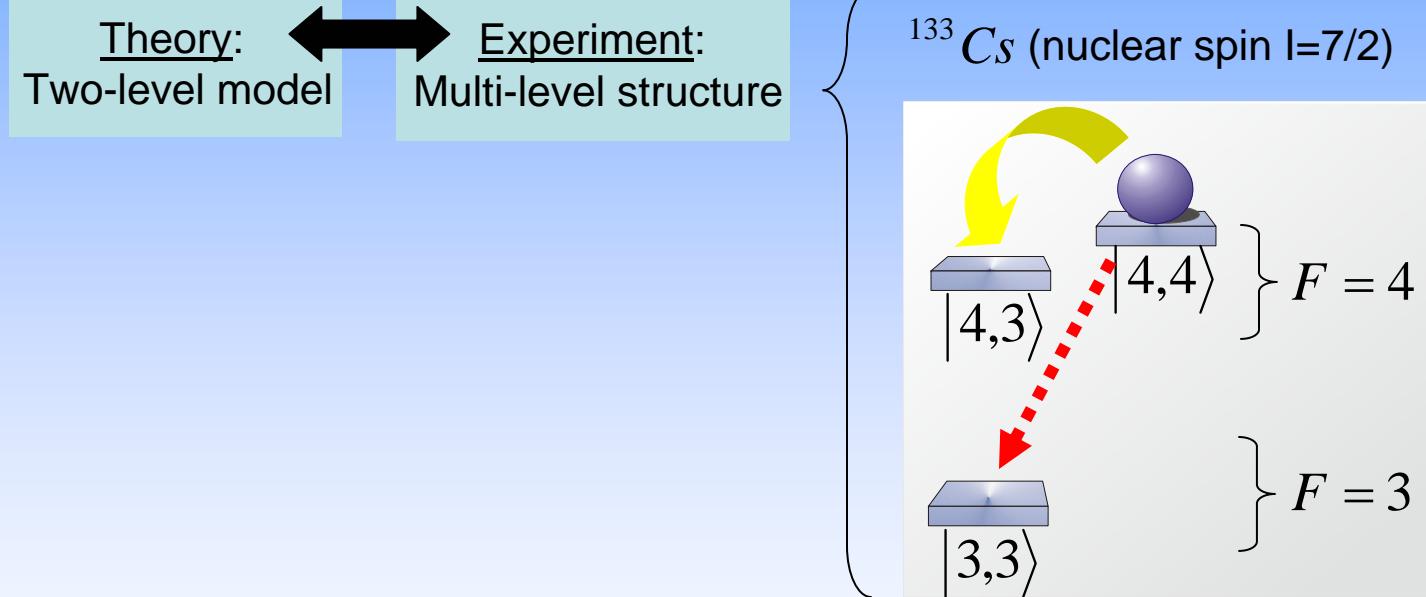


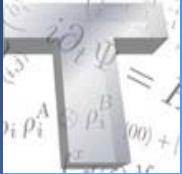
# Experimental realization of purely dissipation based entanglement:

Theory:  $\longleftrightarrow$  Experiment:  
Two-level model                    Multi-level structure



# Experimental realization of purely dissipation based entanglement:





# Experimental realization of purely dissipation based entanglement:

Theory:  
Two-level model

Experiment:  
Multi-level structure

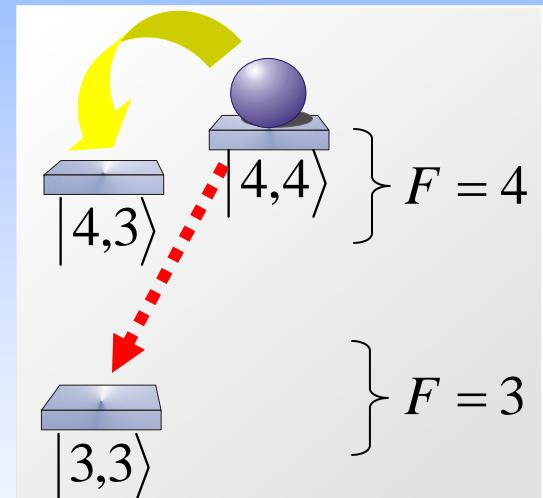
Quasi steady-state:

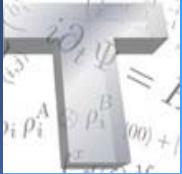
$$\xi = \frac{8}{P_2 + 7} \frac{\tilde{\Gamma} + d\Gamma P_2^2 (\mu - \nu)^2}{\tilde{\Gamma} + dTP_2}$$

$P_2(t)$ : polarization

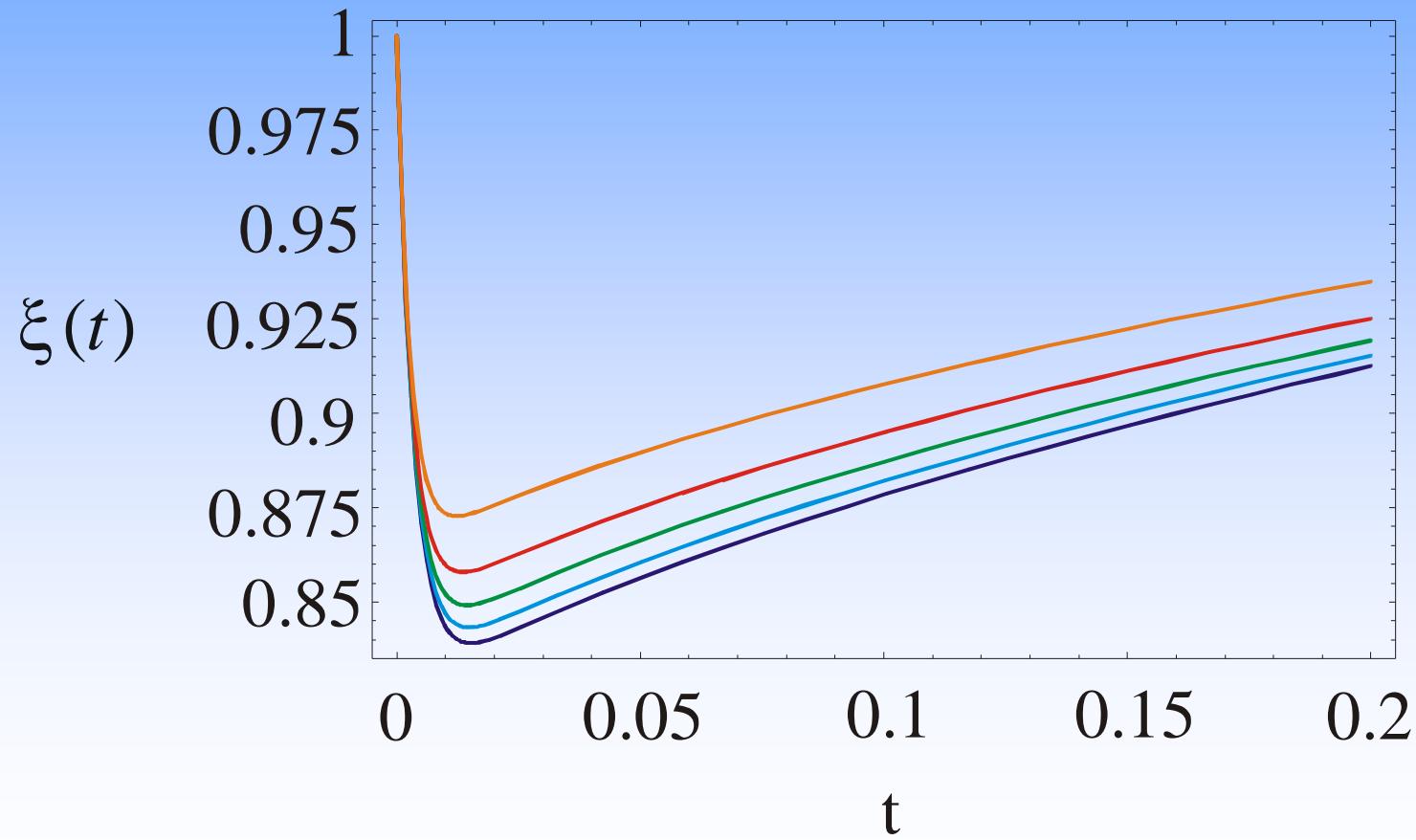
$n_2(t)$  : depopulation

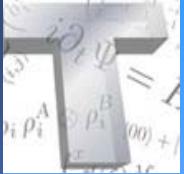
$^{133}Cs$  (nuclear spin  $I=7/2$ )





## Experimental realization of purely dissipation based entanglement:





## Summary:



### **Proposal for long-lived entanglement**

- Atomic ensembles at room temperature
- Dissipatively driven entanglement
- Environment = vacuum modes of the electromagnetic field



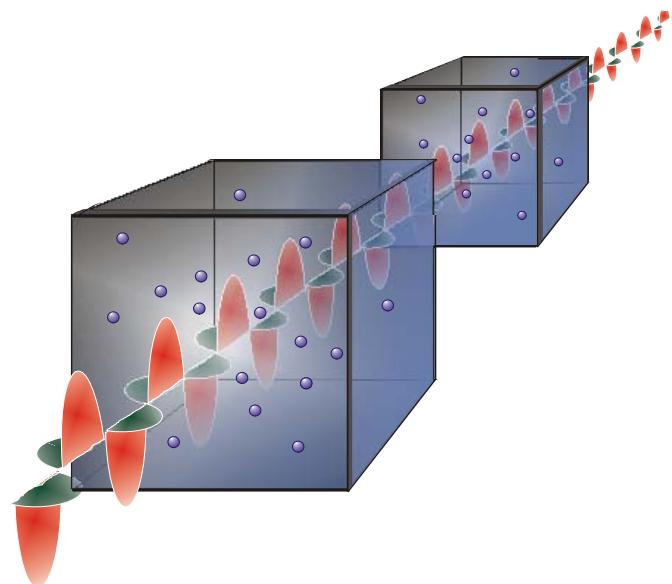
### **Experimental results**

- Quasi steady state



### **Future**

- Multi-level systems
- More systems



A dark red curtain is drawn back on either side of a central black rectangular area, creating a stage-like effect. The background outside the stage area is a solid blue.

*Thank you very  
much for your  
attention*