

Quantum Control in Cavity Optomechanics: Theory and Experiment

Sebastian Hofer

Aspelmeyer group (Vienna) Klemens Hammerer (Hannover)





Cavity optomechanical systems





Micromirrors

Aspelmeyer (Vienna) Heidmann (Paris) Bouwmeester (St Barbara, Leiden)



Micromembranes

Harris (Yale) Kimble (Caltech) Treutlein (Basel)



Microtoroids

Kippenberg (MPQ) Weig (LMU) Vahala (Caltech) Bowen (UQ)





Levitated nanoobjects

Raizen (Austin) Barker (London) Aspelmeyer

+ more massive systems (e.g., GW detectors)
+ electromechanical setups (microwave)



simple model:

optical + mechanical resonator coupled via radiation pressure

Applications





Observed quantum effects



- quantum coherent state transfer
- ground state cooling
- quantum coherent coupling
- ponderomotive squeezing
- back-action noise in position sensing
- optomechanical entanglement
- feedback control within decoherence time

O'Connell et al., Nature 464, 697 (2010) Palomaki, Nature 495, 210 (2013)

Chan, Nature 478, 89 (2011) Teufel, Nature 475, 359 (2011)

Verhagen, Nature 482, 63 (2012)

Safavi-Naeini, arXiv:1302.6179 (2013) Brooks, Nature 488, 476 (2012)

Purdy, Science 339, 801 (2013)

Palomaki, Science 342, 710 (2013)

Wilson, arXiv:1410.6191 (2014)

Challenges + (possible) solutions

Quantum protocols need...

- Overcome coupling to environment
 - passive cooling of bath
 - cut ties to environment by *levitation*
 - environment engineering, shape mode spectrum
 - ...
- Accurate read-out of mechanical state
 - clever data processing: Kalman filtering (CW)
 - *pulsed* read-out: beating the SQL and back-action







 \otimes

- Motivation
 - eliminate clamping
 - high Q: Q \rightarrow 10¹⁰ @ p < 10⁻⁹ mbar
 - low dissipation/decoherence
 - high mass ~ 10⁹ amu
 - control over trap parameters
 - free-fall experiments, matterwave interferometry

MM

Romero-Isart et al., NJP 12, 033015 (2010) Chang et al., PNAS 107, 1005 (2010) P. F. Barker et al., PRA 81, 023826 (2010)

in principle: Quantum experiments @ 300K with macroscopic objects!



- 1d cavity cooling
 - in axial direction
 - room temperature
 - ~4 mbar



Rates comparable to standard cavity optomechanics setups, but...



- What we want: low pressure!
 - Q~25 0 4 mber Q~10° 0 10-7 mber

ground state cooling for p < 10⁻⁷ mbar



• No particle has been trapped in UHV so far!

	particle size	final pressure	cooling
Ashkin et al. APL, 19(8):283, 1971	20µm	10 ⁻⁶ mbar	3d fb cooling
Li, <i>PhD thesis</i> , University of Texas, 2011	1-5µm	10 ⁻⁶ mbar	3d fb cooling
Gieseler et al. <i>PRL 109</i> , 103603, 2012	140nm	10 ⁻⁶ mbar	3d fb cooling
Asenbaum et al. <i>Nature Comm 4</i> , 2743 2013	100nm-1µm	10 ⁻⁸ mbar	1d cavity cooling (no trapping)
Kiesel et al. PNAS 110, 14180-14185, 2013	254nm	1 mbar	1d cavity cooling
Monteiro et al. <i>NJP</i> , 15:015001, 2013.	20-500nm	5 mbar	no
Millen et al. <i>Nature Nano</i> , 9:425–429, 2014	50nm-2.56µm	1 mbar	no
Moore et al. <i>PRL</i> 113, 251801	5µm	10⁻ ⁷ mbar	3d fb cooling

(

particles lost at low pressures



radial direction: thermally activated escape • solution: 3d feedback cooling



Li et al., Nature Physics 7, 527 (2011) Gieseler et al., PRL 109, 103603 (2012)

Real-time optimal state estimation



 \bigotimes

Real-time optimal state estimation

- Goal: find optimal estimate of full quantum state
- CW measurement
- reconstruction of mechanical state from measuring light
- Kalman filter
 - Gaussian systems: KF = optimal estimator (minimum-mean-square error)
 - real-time estimator \rightarrow feedback
 - based on dynamical model of system
- Quantum filter
 - *Gaussian* systems: KF solves the *stochastic master equation* (for homodyne detection)
 - obtain conditional quantum state

KF for mechanical systems: Finn et al., PRD 63, 062004 (2001) Iwasawa et al., PRL 111, 163602 (2013)

KF for quantum systems: Yonezawa et al., Science 337:1514 (2012) Geremia et al., PRL 91, 250801 (2003)

V.P. Belavkin (1980)

R.E. Kalman (1960)





Kalman filter in a nutshell



• algorithm

0) initial state

- 1) propagation = prediction
- 2) measurement update (Bayesian conditioning)



Kalman filter Vienna system



- requires accurate system and measurement model
- must account for:
 - thermal noise
 - technical noise sources
 - broadband, colored laser noise
 - narrowband noise peaks
 - multimode structures
- model must be validated against measurements







Kalman filter Vienna system

 \bigotimes





W. Wieczorek, SH et al., submitted (2015)

Beating the SQL by QND

- Pulsed state state tomography
 - stroboscopic measurement of $\, x_m \,$
 - measurement back-action on $\,p_m\,$
 - operated on resonance, for a bad cavity



- "Cooling by measurement"
 - = conditional reduction of variance
 - interaction strength (no cavity) $\chi\approx 1/\sigma_x\approx 10^{-4}$



M. Vanner et al., Nature Communications 4 (2013)



Beating the SQL by QND



- Goal: $\chi > 1 \Leftrightarrow \sigma_x < 1$
 - conditionally squeezed state
- How?
 - large single-photon coupling
 - interaction enhanced by cavity
 - $\chi \approx 4$
- Problem:
 - diffusion due to multiple modes!





Beating the SQL by QND



subtract predictions from measurements



Hamiltonian



optomechanical cavity



$$\kappa = \frac{c\pi}{2LF}$$
$$\bar{n} \simeq \frac{k_B T}{\hbar \omega_m}$$
$$\Delta = \omega_c - \omega_f$$

 linearised interaction Hamiltonian (for strong driving)



 $-\omega_m$

 ω_{l}

 ω_{c}

$$H_{rp} = g(a_c + a_c^{\dagger})(a_m + a_m^{\dagger})$$

$$= g(a_c a_m^{\dagger} + a_c^{\dagger} a_m) + g(a_c a_m + a_c^{\dagger} a_m^{\dagger})$$

$$\int \int \int d_{\kappa} d_{\kappa}$$

Steady-state phase diagram



Conditional-state phase diagram

• conditioned on homodyne detection of phase quadrature



accepted in PRA, arXiv:1411.1337

7

Feedback cooling (homodyne detection)



 \bigotimes

Feedback cooling with a twist





Feedback cooling with a twist





Time-continuous teleportation



 \bigotimes

CV Bell measurement/teleportation



- CV Bell measurement = 2 homodyne detectors + beam splitter
 - measures EPR quadratures $X_{-} = X_{A} X_{B}$ $P_{+} = P_{A} + P_{B}$ with outcomes m, and m,
- CV teleportation
 - Bell measurement projects mechanics onto into displaced input state
 - displacement by m_x and m_p recovers input state

Hofer et al. PRA 84, 052327, (2011) Palomaki et al., Science 342, 710 (2013)

- continuous operation of teleportation CW, no pulses
- dissipative remote state preparation

$$\dot{\rho} = \mathcal{D}[J]\rho \quad \underbrace{t \to \infty}_{J|\psi_{\rm in}} \rho_{\rm ss} = |\psi_{\rm in}\rangle\langle\psi_{\rm in}|$$

teleportation of a squeezed state





Time-continuous teleportation





Time-continuous ent. swapping



 (\mathbb{K})

Time-continuous ent. swapping

- create steady-state entanglement between two mechanical oscillators
- dissipative remote state preparation

$$\dot{\rho} = \mathcal{D}[J]\rho \quad \underbrace{t \to \infty}_{J|\psi_{\text{EPR}}} \approx \underbrace{t \to \infty}_{0} \quad \rho_{\text{ss}} \approx |\psi_{\text{EPR}}\rangle \langle \psi_{\text{EPR}}|$$











Acknowledgements

The Aspelmeyer Group

Thank you!

Klemens

Hammerer

@ | Godany