

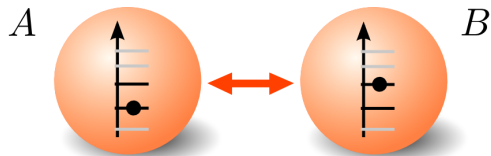
# Photonic simulation of quantum systems: combating noise and weak interaction

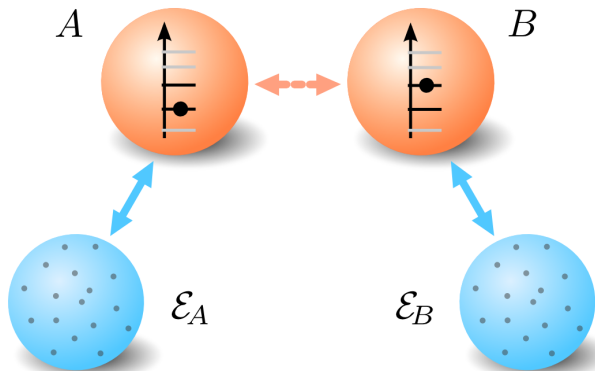
**Michal Mičuda, Martina Miková, Ivo Straka, Robert Stárek,  
Miroslav Ježek, Miloslav Dušek, Radim Filip, Jaromír Fiurášek**



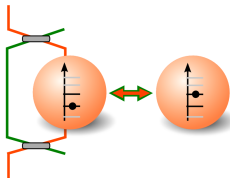
Department of Optics, Faculty of Science  
Palacký University Olomouc

Photons Beyond Qubits Workshop, Mar 18 2015

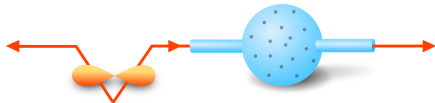




- Quantum C-Z gate for weakly interacting qubits
  - boosting interaction by interference
  - photonic simulator



- Entanglement transfer through noisy environment
  - multi-particle incoherent environment
  - environment probing
  - photonic simulators



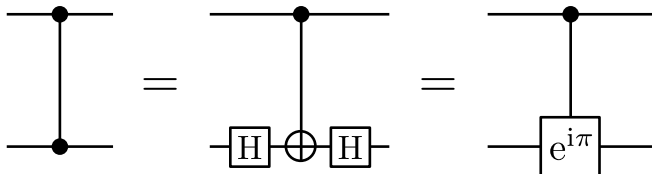
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- 2 Entanglement transfer through noisy environment

$$CZ = I - 2|11\rangle\langle 11|$$

$$CZ|jk\rangle = (-1)^{jk}|jk\rangle$$

$$CPHASE = \exp(i\phi|11\rangle\langle 11|)$$

$$\phi = \pi$$



spin-spin coupling

# Quantum C-Z gate of weakly spin-spin coupled qubits

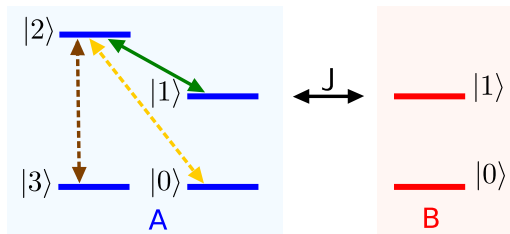
$$U_\phi = \exp(i\phi |11\rangle\langle 11|)$$

$$\phi < \pi$$

Auxiliary BS coupling  $|1\rangle \rightleftharpoons |2\rangle$

$$|1\rangle \rightarrow t|1\rangle + r|2\rangle$$

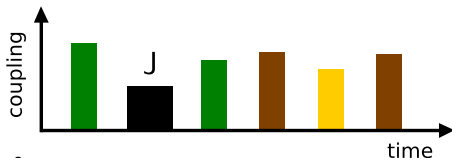
$$|2\rangle \rightarrow \tilde{t}^*|2\rangle - \tilde{r}^*|1\rangle$$



Projection back to  $\{|0\rangle, |1\rangle\}$

$$(t\tilde{t} - r\tilde{r}^*)|0\rangle\langle 0| + |1\rangle\langle 1|$$

$$\text{CZ: } r\tilde{r}^*/t\tilde{t} = (1 + e^{i\phi})/2$$



$$\text{Probability of success: } P_s = |t\tilde{t} - r\tilde{r}^*|^2$$

$$\text{maximized for } |t|^2 = |\tilde{t}|^2 = 1/[1 + |\cos(\phi/2)|]$$

# Quantum C-Z gate of BS-coupled qubits

Beam splitter coupling of two bosonic modes

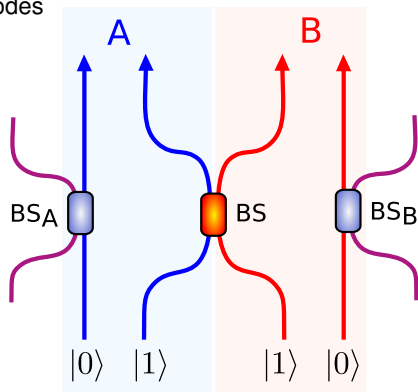
$$U_{BS}(t) = \exp[\kappa t(ab^\dagger - a^\dagger b)]$$

$$U_{BS}(R) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \sqrt{1-R} & 0 & 0 \\ 0 & 0 & \sqrt{1-R} & 0 \\ 0 & 0 & 0 & 1-2R \end{pmatrix}$$

Strong interaction

$$R = \sin^2(\kappa t) > 1/2$$

$$U_{BS}(2/3) \sim CZ$$





# Quantum C-Z gate of weakly BS-coupled qubits

Weakly coupled systems

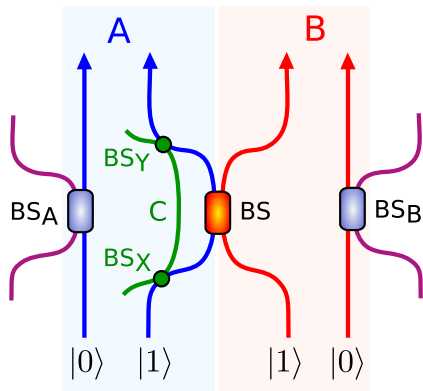
$$R = \sin^2(\kappa t) < 1/2$$

BS partially bypassed by C

$BS_{X,Y}$ : amp. transmittances  $t_X, t_Y$

$$\text{C-Z: } \frac{r_X r_Y}{t_X t_Y} = \frac{3R - 2}{2\sqrt{1 - R}}$$

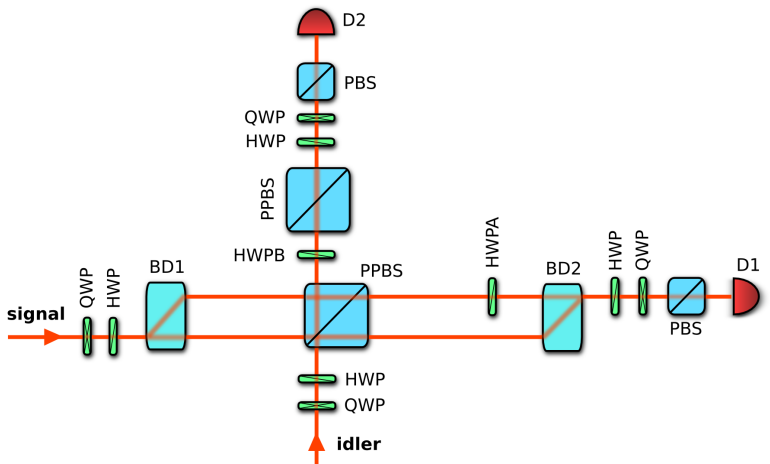
Probability of success:  $R^2 \frac{T_X T_Y}{4}$   
 maximized for  $T_X = T_Y$



$$U = \begin{pmatrix} t_{AtB} & 0 & 0 & 0 \\ 0 & t_{At} & 0 & 0 \\ 0 & 0 & t_B(t t_X t_Y + r_X r_Y) & 0 \\ 0 & 0 & 0 & (2t^2 - 1)t_X t_Y + t r_X r_Y \end{pmatrix} = R \frac{t_X t_Y}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

# Photonic simulator of C-Z gate with weak coupling and bypass

BS: PPBS  $R_V = 1/3$ ,  $R_H = 0$

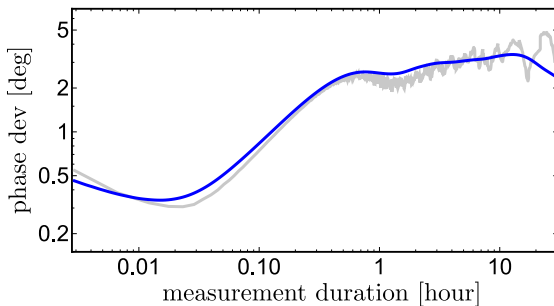
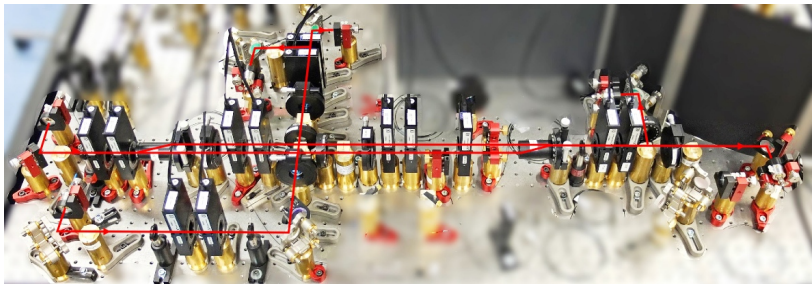




# Photonic simulator of C-Z gate with weak coupling and bypass

Experimental setup based on Toffoli gate experiment

[M. Mičuda et al., Phys. Rev. Lett. 111, 160407 (2013)]

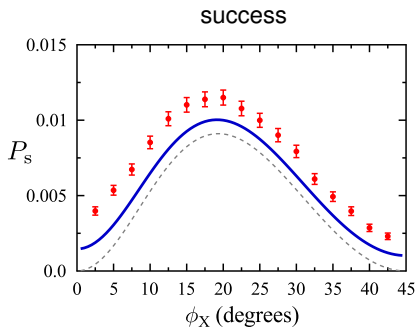
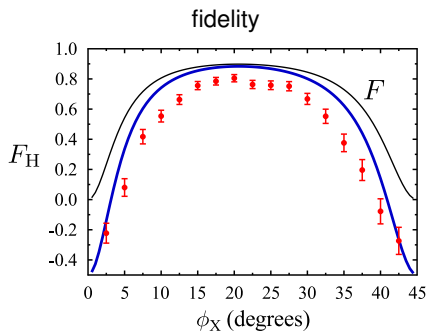


# Simulation results

Gate fidelity:  $F_x = \text{Tr}[\chi\chi_{cz}]/(\text{Tr}[\chi]\text{Tr}[\chi_{cz}])$

Hofmann bound on quantum process fidelity:  $F_x \geq F_1 + F_2 - 1$

[H.F. Hofmann, PRL 94, 160504 (2005); M. Mićuda et al., PRA 89, 042304 (2014)]



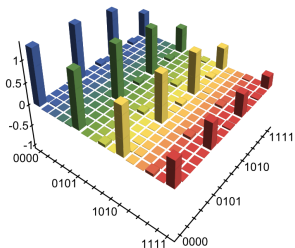
Theoretical model:  $R_H = 1.9\%$ ,  $\mathcal{V}_{\text{HOM}} = 94\%$

Maximum predicted gate fidelity: 88.9% at  $\phi_X = 20^\circ$

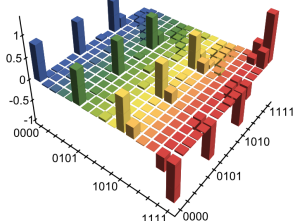
# Simulation results: full tomography of quantum process

Number of preparation/measurement configurations:  $6^4 = 1296$   
10 s acquisition per configuration  $\rightarrow$  total time of 6 hours

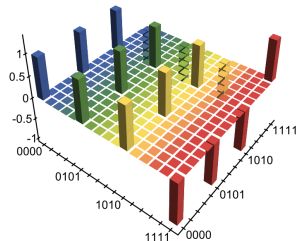
without bypass



optimal bypass



ideal C-Z gate



$$F_{\chi} = 84.6\%$$

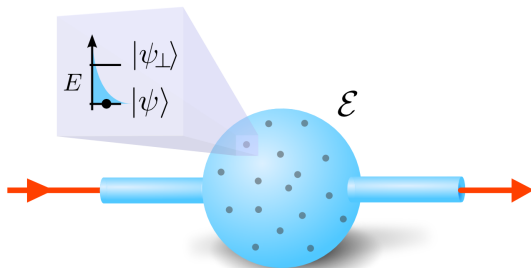
[M. Mičuda, R. Stárek, I. Straka, M. Miková, M. Dušek, M. Ježek, R. Filip, and J. Fiurášek, submitted 2015]

- 1 Quantum C-Z gate for weakly interacting qubits
- 2 Entanglement transfer through noisy environment

Multi-qubit environment  $\mathcal{E}$ :

$$\rho = (1 - p_T)|\psi\rangle\langle\psi| + p_T|\psi_\perp\rangle\langle\psi_\perp|$$

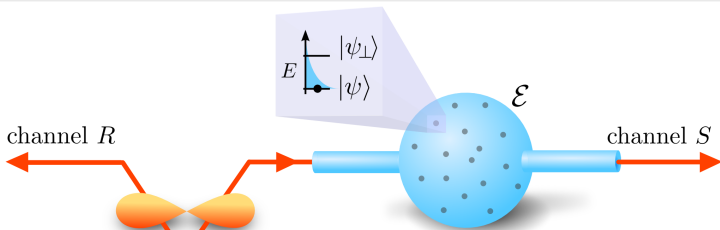
$$p_T = \frac{\exp\left(-\frac{\Delta E}{k_B T}\right)}{1 + \exp\left(-\frac{\Delta E}{k_B T}\right)}$$



Incoherent environment: qubits do not interfere or interact with another qubits



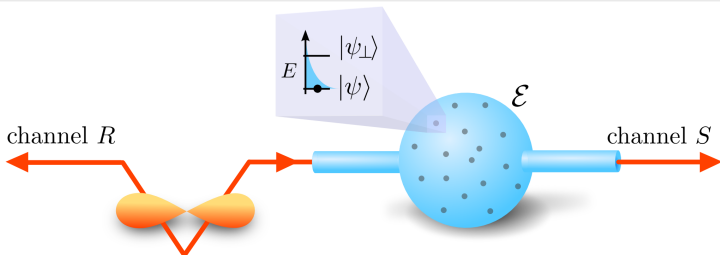
# Entanglement transfer through incoherent environment



$$\text{Input state } |\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\psi\rangle|\psi_\perp\rangle - |\psi_\perp\rangle|\psi\rangle)$$

$$\text{Output state } \rho_{RS} = P_S |\Psi^-\rangle_{RS} \langle\Psi^-| + (1 - P_S) \frac{1_R}{2} \otimes \rho_S$$

# Entanglement transfer through incoherent environment



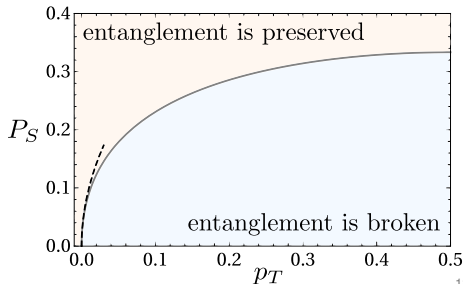
Input state  $|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|\psi\rangle|\psi_\perp\rangle - |\psi_\perp\rangle|\psi\rangle)$

Output state  $\rho_{RS} = P_S |\Psi^-\rangle_{RS} \langle\Psi^-| + (1 - P_S) \frac{1_R}{2} \otimes \rho_S$

The state remains entangled if

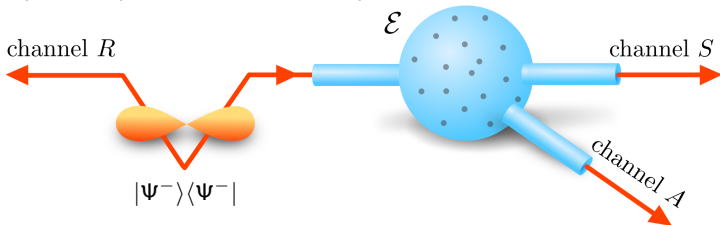
$$P_S > \frac{\sqrt{p_T(1-p_T)}}{1 + \sqrt{p_T(1-p_T)}}$$

$$\frac{p_T}{P_S^2} < 1 \text{ for } p_T \ll 1$$



# Using additional environment probing

Three possible processes: success, flip, and error;  $P_S + P_F + P_E = 1$

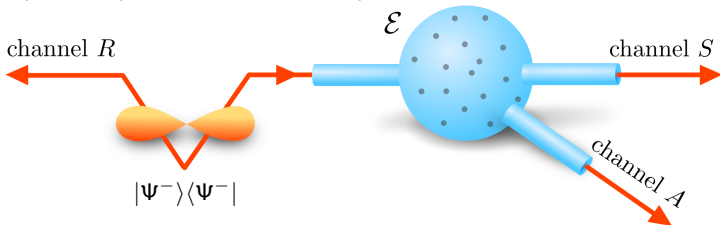


Projecting channel A to environment's ground state gives

$$\rho_{RS} \propto (1 - p_T)P_S|\Psi^-\rangle_{RS}\langle\Psi^-| + \frac{P_F}{2}|\psi_\perp\rangle_R\langle\psi_\perp| \otimes \rho_S + (1 - p_T)P_E\frac{1_R}{2} \otimes \rho_S$$

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The state remains entangled if

$$P_S > \frac{1}{2} \left( \sqrt{p_T P_E (4 - 3p_T P_E)} - p_T P_E \right)$$

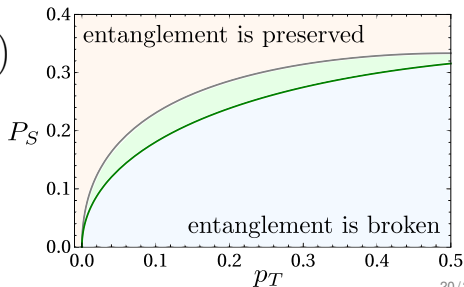
$$\frac{p_T P_E}{P_S^2} < 1 \text{ for } p_T \ll 1$$

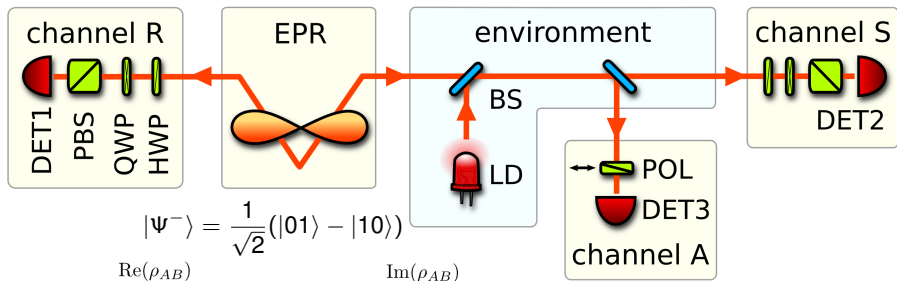
$P_E = 0$  for single particle environment

[E. Nagali et al., PRA 79, 060304(R) (2009)]

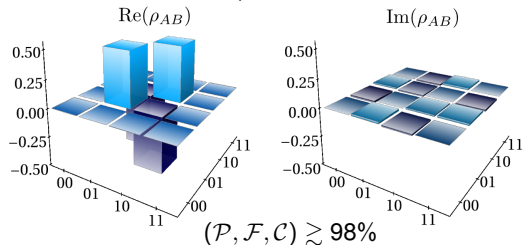
[M. Gavenda et al., PRA 81, 022313 (2010)]

[M. Gavenda et al., PRA 83, 042320 (2011)]





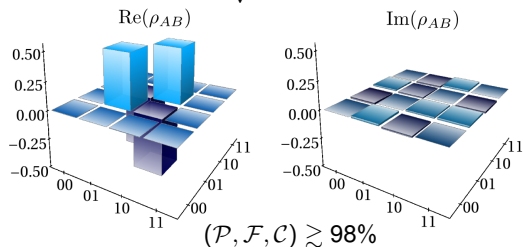
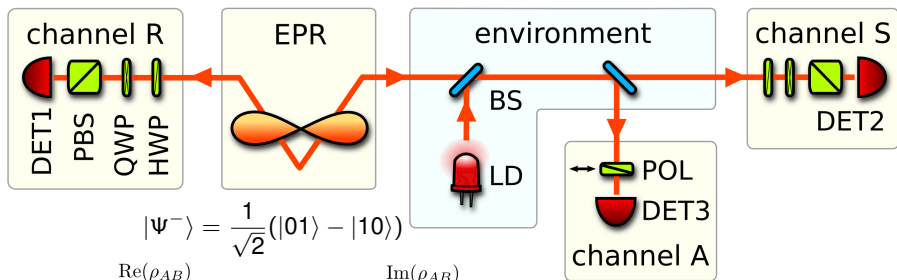
$$|\psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$



[M. Miková et al., PRA 90, 022317 (2014)]

$|0\rangle\langle 0|$  conditioning (or Tr)

- 3-fold coincidences
- Full tomography
- Coincidence window
- Noise depolarization



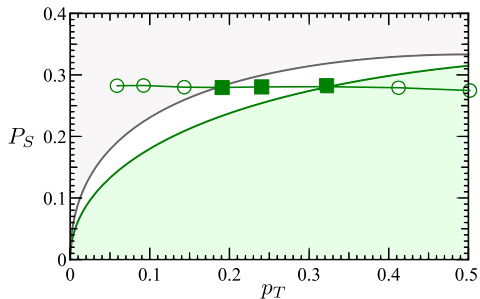
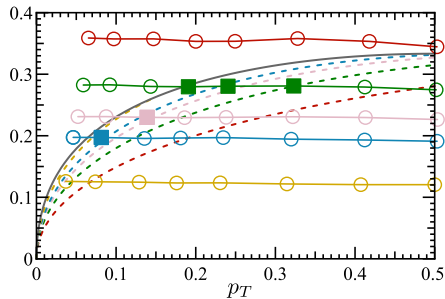
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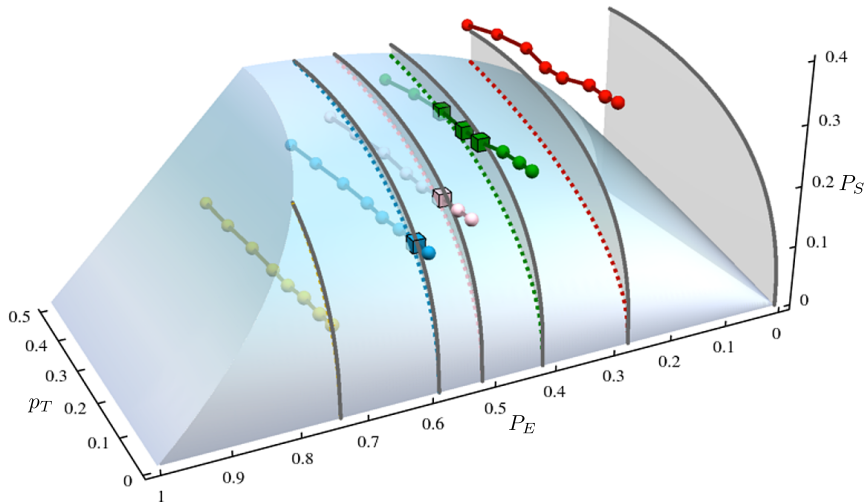
Model of the simulator:

$$\rho_{RS} \propto (1 - p_T)|\psi^-\rangle_{RS}\langle\psi^-| + \frac{1}{2}|1\rangle\langle 1| \otimes \rho_S + (1 - p_T)\tilde{P}_E \frac{1}{2} \otimes \rho_S, \quad \tilde{P}_E \propto \frac{\tau R_S R_N}{R_{\psi^-}}$$

$P_E = 0.42$  $P_E = 0.27, 0.42, 0.52, 0.58, 0.74$ 

# Simulation results and generalization of the simulator

$$\text{Model: } \rho_{RS} \propto (1 - p_T) |\Psi^-\rangle_{RS} \langle \Psi^-| + \frac{1}{2} |1\rangle_R \langle 1| \otimes \rho_S + (1 - p_T) \tilde{P}_E \frac{1_R}{2} \otimes \rho_S$$

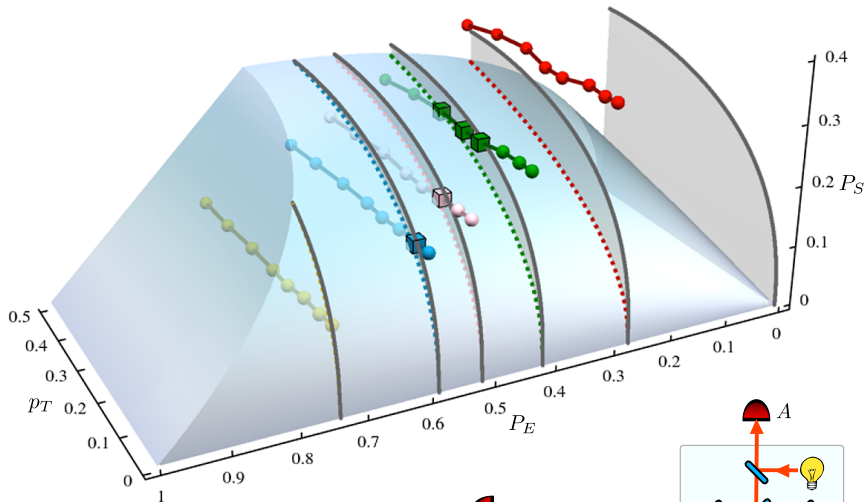


[I. Straka, M. Miková et al., in preparation 2015]

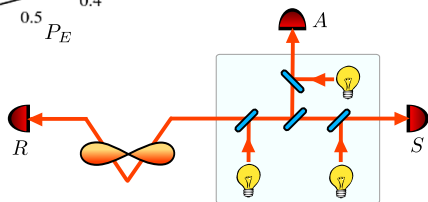


# Simulation results and generalization of the simulator

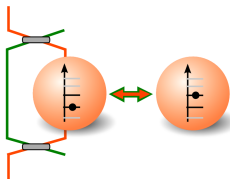
$$\text{Model: } \rho_{RS} \propto (1 - p_T) |\Psi^-\rangle_{RS} \langle \Psi^-| + \frac{1}{2} |1\rangle_R \langle 1| \otimes \rho_S + (1 - p_T) \tilde{P}_E \frac{1}{2} \otimes \rho_S$$



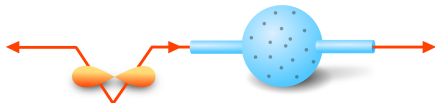
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Thank you for your attention!



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