

Distributing Entanglement with Separable States

Ladislav Mišta,¹ Christian Peuntinger,² Vanessa Chille,² Natalia Korolkova,³ Michael Förtsch,² Jan Korger,² Christoph Marquardt,² and Gerd Leuchs²

¹Department of Optics, Palacký University, 17. listopadu 12, 771 46 Olomouc, Czech Republic

²Max Planck Institute for the Science of Light, Günther-Scharowsky-Str. 1/Bldg. 24, Erlangen, Germany

³School of Physics and Astronomy, University of St. Andrews, North Haugh, St. Andrews, Fife, KY16 9SS, Scotland

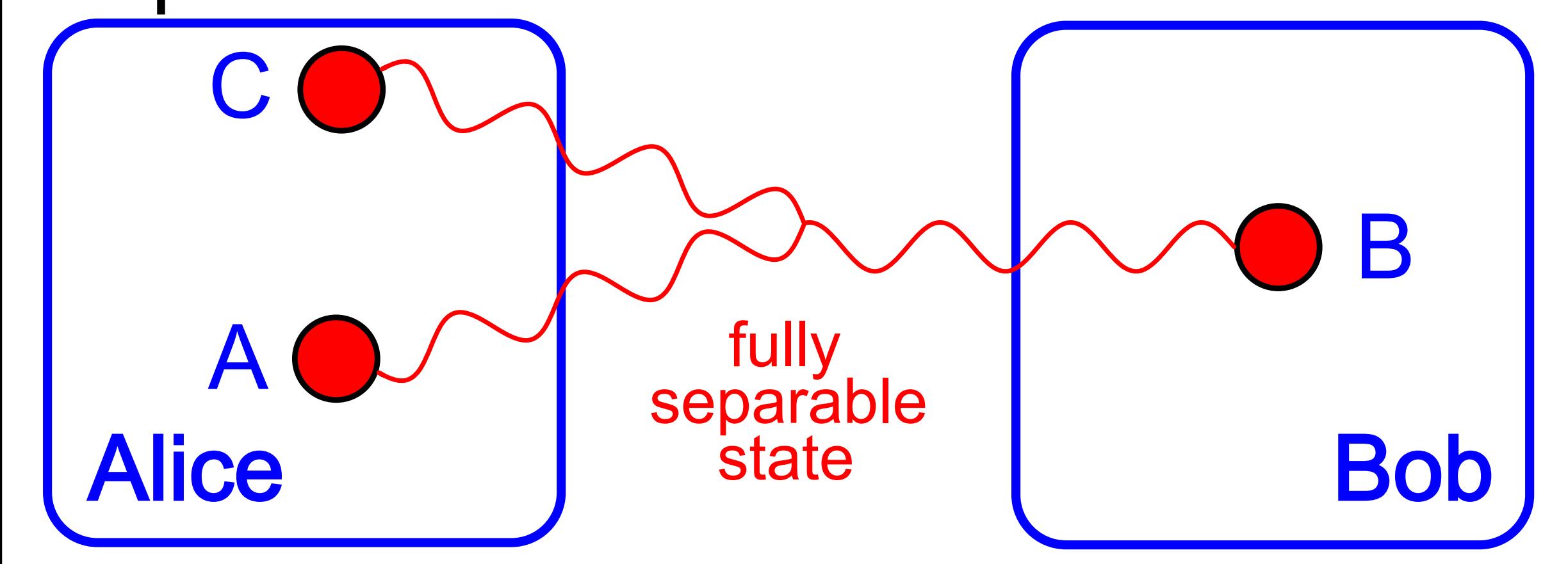
Entanglement distribution by separable states

Entanglement distribution without sending entanglement [1].

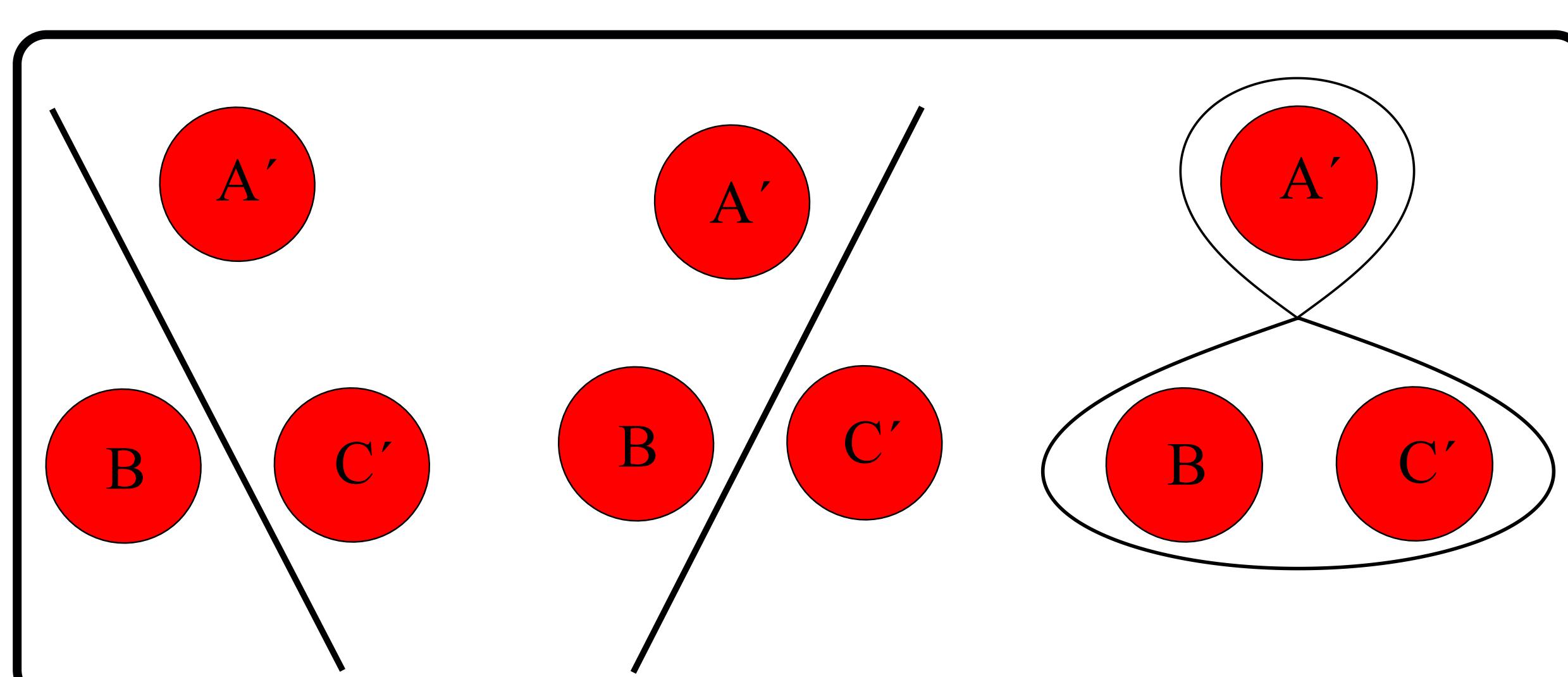
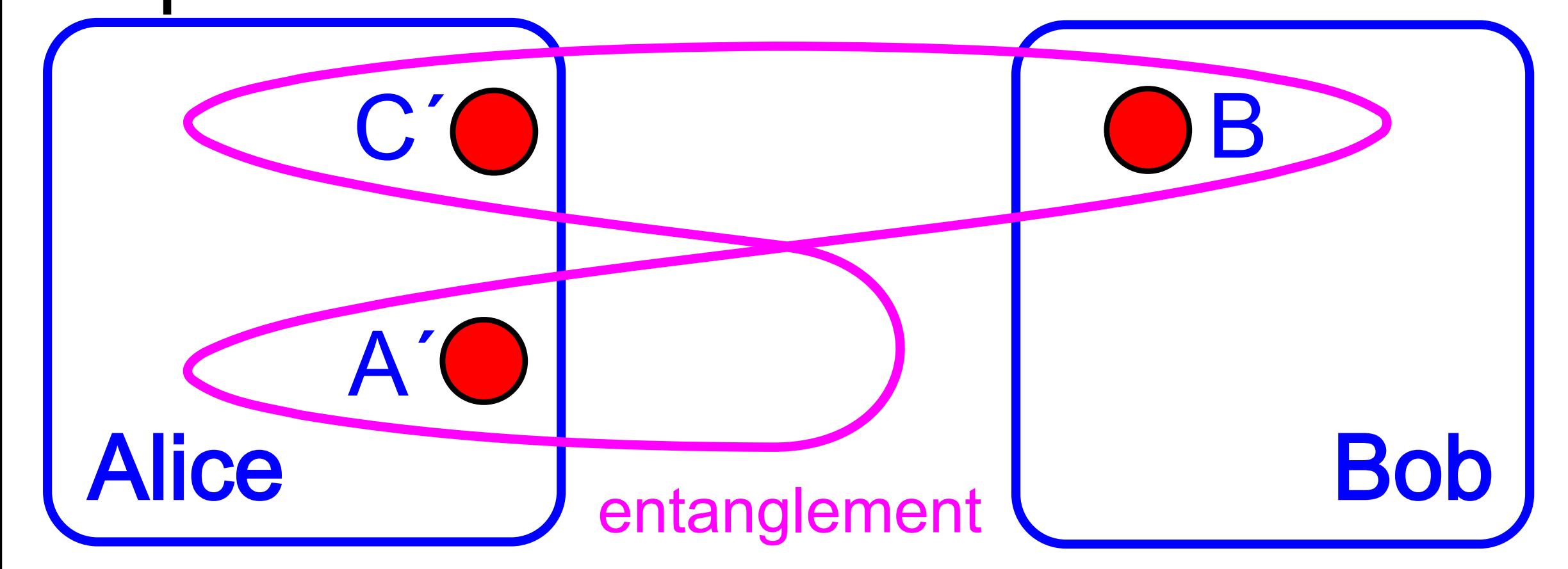
Why is it interesting?

1. Application of multipartite bound entanglement.
2. Construction of new cryptographic concepts [2, 3].
3. Interpretation of quantum discord [4].

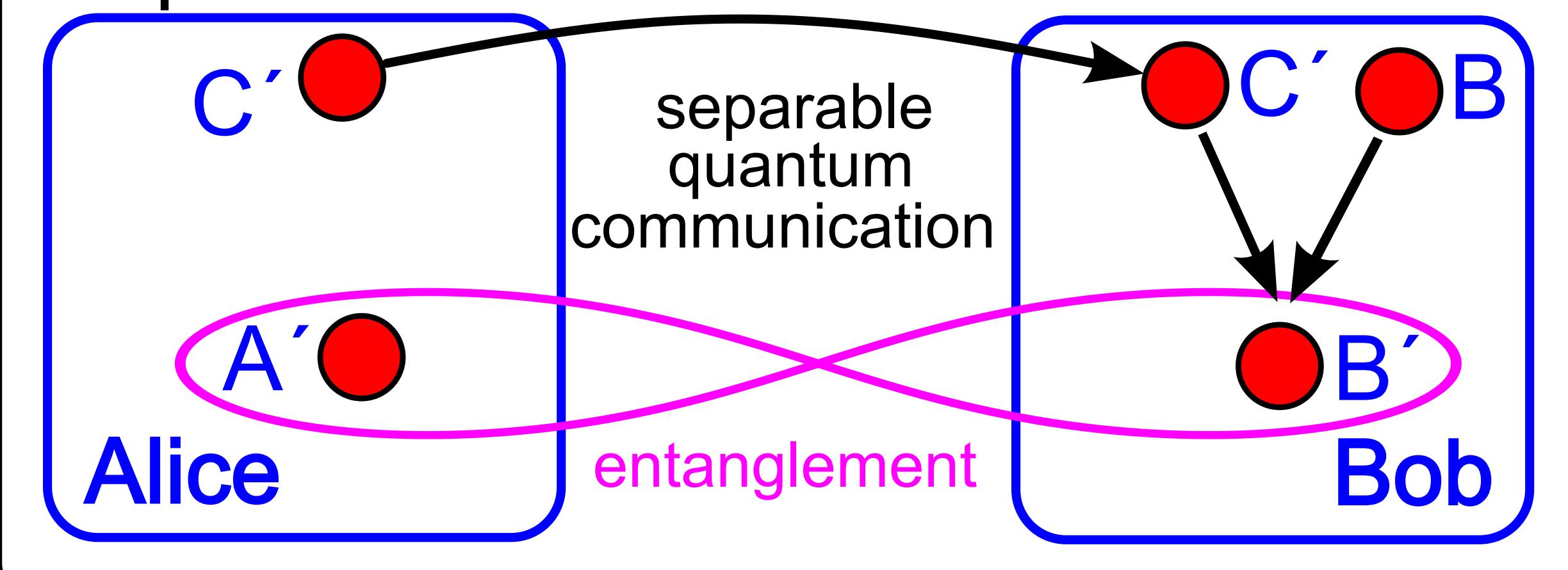
Step 1



Step 2



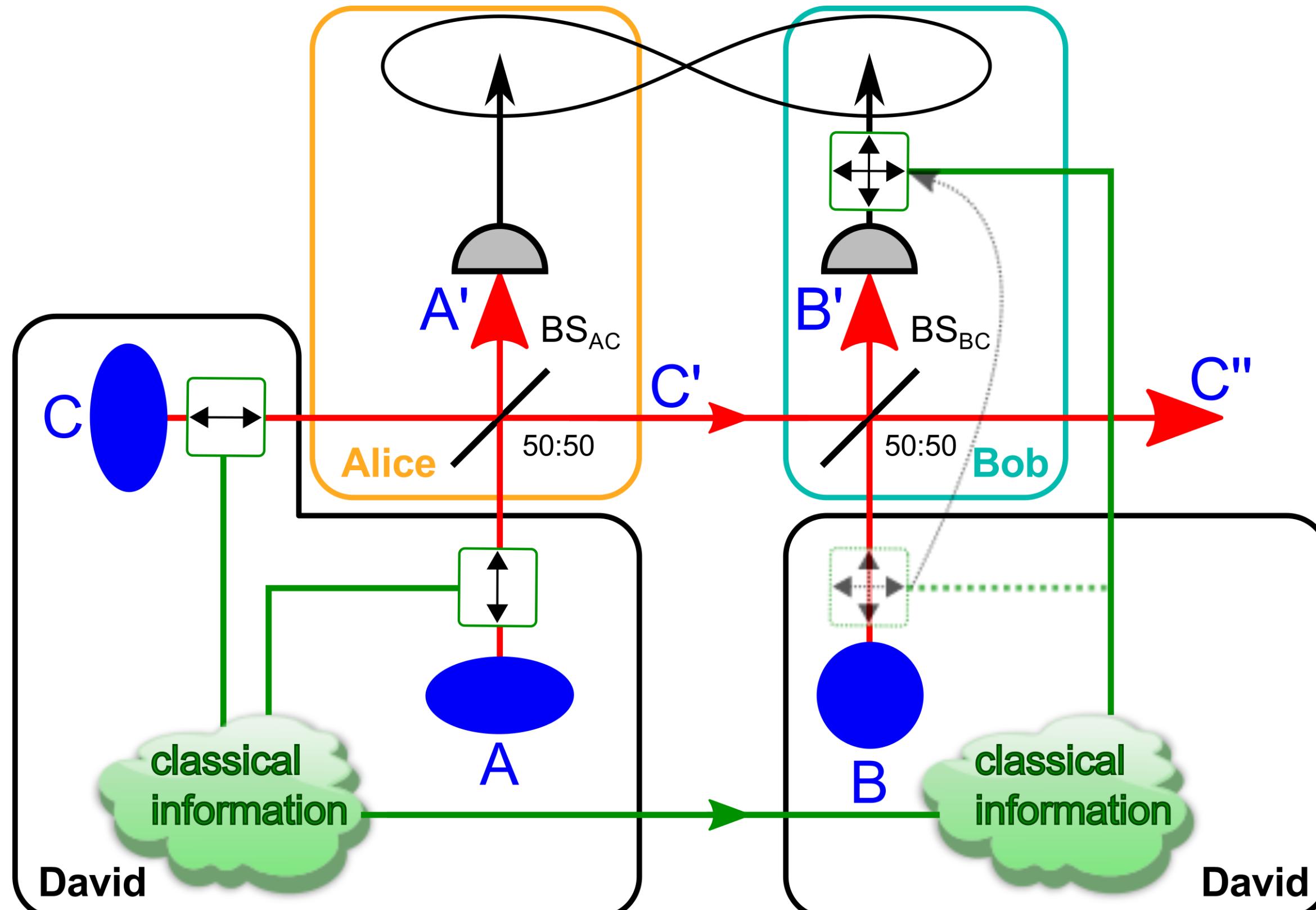
Step 3



References

- [1] T. S. Cubitt, F. Verstraete, W. Dür, and J. I. Cirac, Phys. Rev. Lett. **91**, 037902 (2003).
- [2] J. Bae, T. Cubitt, and A. Acin, Phys. Rev. A **79**, 032304 (2009).
- [3] L. Mišta, Jr. and N. Korolkova, Phys. Rev. A **86**, 040305 (2012).
- [4] A. Streltsov, H. Kampermann, and D. Bruß, Phys. Rev. Lett. **108**, 250501 (2012); T. K. Chuan, J. Maillard, K. Modi, T. Paterek, M. Paternostro, and M. Piani, Phys. Rev. Lett. **109**, 070501 (2012); A. Kay, Phys. Rev. Lett. **109**, 080503 (2012).
- [5] G. Giedke, B. Kraus, M. Lewenstein, and J. I. Cirac, Phys. Rev. A **64**, 052303 (2001).
- [6] V. Giovannetti, S. Mancini, D. Vitali, and P. Tombesi, Phys. Rev. A **67**, 022320 (2003).
- [7] C. E. Vollmer, D. Schulze, T. Eberle, V. Händchen, J. Fiurášek, and R. Schnabel, Phys. Rev. Lett. **111**, 230505 (2013).
- [8] A. Fedrizzi, M. Zuppardo, G. G. Gillett, M. A. Broome, M. de Almeida, M. Paternostro, A. G. White, and T. Paterek, Phys. Rev. Lett. **111**, 230504 (2013).

Scheme with squeezed light



Step 1: David prepares a three-mode fully separable state. Modes A and C are squeezed, mode B is vacuum,

$$\hat{x}_{A,C} = e^{\pm r} \hat{x}_{A,C}^{(0)}, \quad \hat{p}_{A,C} = e^{\mp r} \hat{p}_{A,C}^{(0)}, \quad \hat{x}_B = \hat{x}_B^{(0)}, \quad \hat{p}_B = \hat{p}_B^{(0)},$$

and all modes are displaced by x and p , $\langle x^2 \rangle = \langle p^2 \rangle = \frac{e^{2r}-1}{2}$, as

$$\begin{aligned} \hat{p}_A &\rightarrow \hat{p}_A - p, & \hat{x}_C &\rightarrow \hat{x}_C + x, \\ \hat{x}_B &\rightarrow \hat{x}_B + \sqrt{2}x, & \hat{p}_B &\rightarrow \hat{p}_B + \sqrt{2}p. \end{aligned}$$

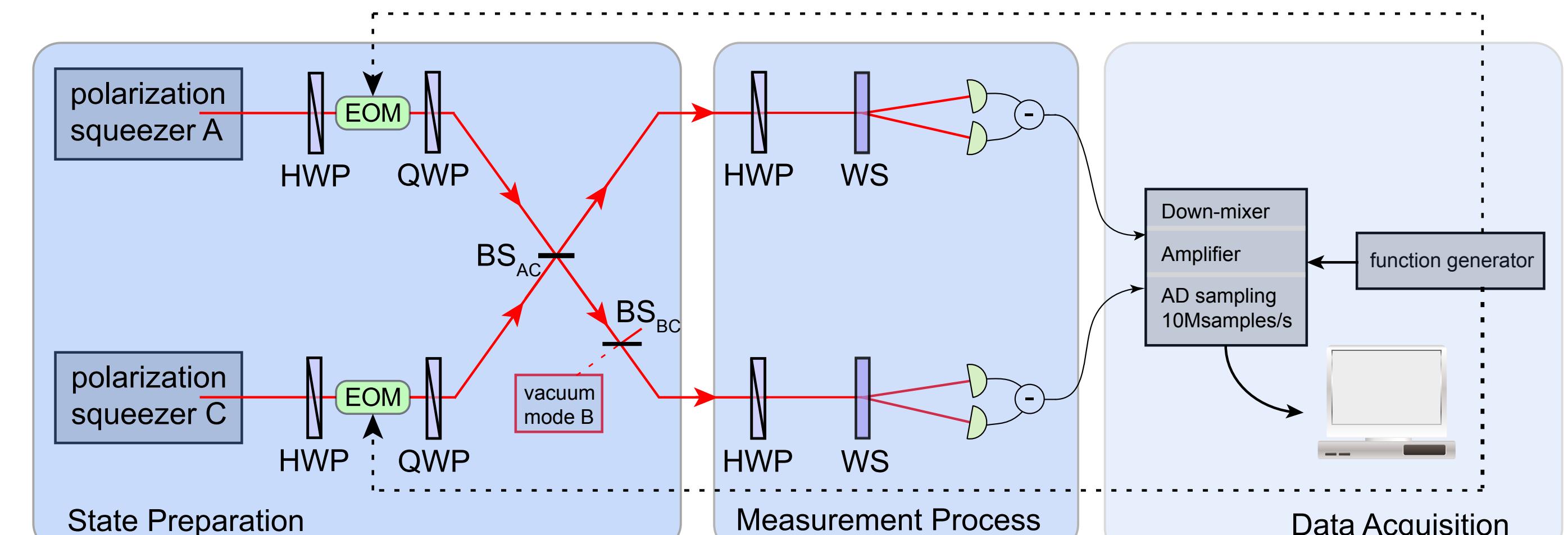
Step 2: Alice creates by a beam splitter BS_{AC} entanglement between A' and (BC') and both C' remains separable from $(A'B)$ and B from $(A'C')$.

Alice sends a separable mode C to Bob.

Step 3: Bob entangles A' with B' by a beam splitter BS_{BC} .

L. Mišta, Jr. and N. Korolkova, Phys. Rev. A **80**, 032310 (2009).

Experiment



Quadratures: $\hat{x} = \frac{\hat{S}_1}{\sqrt{2S_3}}$, $\hat{p} = \frac{\hat{S}_2}{\sqrt{2S_3}}$, $\hat{S}_3 \approx S_3$, $\hat{S}_{1,2,3}$ —Stokes operators.

Measurement: covariance matrix γ of modes A' , B and C' .

Separability certification in step 2 (PPT criterion [5]):

mode j is separable $\Leftrightarrow \gamma^{(T_j)} + i\Omega \geq 0$,

$$\gamma^{(T_j)} = \Lambda_j \gamma \Lambda_j, \quad \Lambda_j = \sigma_z^{(j)} \bigoplus_{k \neq j} I^{(k)}, \quad \Omega = \bigoplus_{i=1}^3 i \sigma_y.$$

j	A'	B	C'
$\min[\text{eig}(\gamma^{(T_j)} + i\Omega)] \times 10$	-1.44 ± 0.01	5.28 ± 0.03	3.51 ± 0.02

Entanglement recovery in step 3 (product criterion [6]):

$$D \equiv \Delta_{\text{norm}}^2(g\hat{x}_{A'} + \hat{x}_{B'}) \Delta_{\text{norm}}^2(g\hat{p}_{A'} - \hat{p}_{B'}) < 1 \Rightarrow A' - B' \text{ entanglement}.$$

$$D = 0.6922 \pm 0.0002 < 1 \text{ for } g_{\text{opt}} = 0.4235 \pm 0.0005$$

Ch. Peuntinger, V. Chille, L. Mišta, Jr., N. Korolkova, M. Förtsch, J. Korger, Ch. Marquardt, and G. Leuchs, Phys. Rev. Lett. **111**, 230504 (2013). Two more experiments [7, 8].